

3. Lattice Model of a Fluid

Lattice Models

Lattice models provide a minimalist, or coarse-grained, framework for describing the translational, rotational, and conformational degrees of freedom of molecules, and are particularly useful for problems in which entropy of mixing, configurational entropy, or excluded volume are key variables. The lattice forms a basis for enumerating different configurations of the system, or microstates. Each of these microstates may have a different energy, which is then used to calculate a partition function.

$$Q = \sum_i e^{-E_i/k_B T} \quad (1)$$

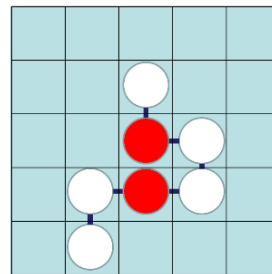
The thermodynamic quantities then emerge from

$$\begin{aligned} A &= -k_B T \ln Q \\ S &= -k_B \sum_i P_i \ln P_i \\ U &= \sum_i P_i E_i \end{aligned}$$

and other internal variables (X) can be statistically described from

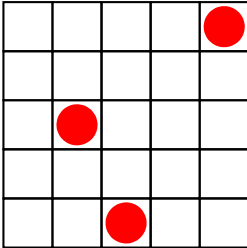
$$\langle X \rangle = \sum_{i=1}^N P_i X_i \quad P_i(E_i) = \frac{e^{-E_i/k_B T}}{Q}$$

We will typically work with a macroscopic volume broken into cells, typically of a molecular size, which we can fill with the fundamental building blocks in our problem (atoms, molecules, functional groups) subject to certain constraints. In this section we will concern ourselves with the mixing of rigid particles, i.e., translational degrees of freedom. More generally, lattice models can include translational, rotational, and conformational degrees of freedom of molecules.



Ideal Lattice Gas

The description of a weakly interacting fluid, gas, solution, or mixture is dominated by the translational entropy or entropy of mixing. In this case, we are dealing with how molecules occupy a volume, which leads to a translational partition function. We begin by defining a lattice and the molecules that fill that lattice:



Parameters:

Total volume: V

Cell volume: v

Number of sites: $M = V/v$

Number of particles: N ($N \leq M$)

Number of contacts each cell has with adjacent cells: z

We begin by assuming that all microstates (configurations of occupied sites in the volume) are equally probable, i.e., $E_i = \text{constant}$. This is the microcanonical ensemble, so the entropy of the fluid is given by Boltzmann's equation

$$S = k_B \ln \Omega \quad (2)$$

where Ω is the number of microstates available to the system. If M is not equal to N , then the permutations for putting N indistinguishable particles into M sites is given by the binomial distribution:

$$\Omega = \frac{M!}{N!(M-N)!} \quad (3)$$

Particles are indistinguishable

Vacancies are indistinguishable

Also, on cubic lattice, we have 6 contacts that each cell makes with its neighbors. The contact number is z , which will vary for 2D ($z = 4$) and 3D ($z = 6$) problems.

How do we choose the size of v ? It has to be considered on a case-by-case basis. The objective of these models is to treat the cell as the volume that a particle excludes to occupation by other particles. This need not correspond to an actual molecular dimension in the atomic sense. In the case of the traditional derivation of the translational partition function for an ideal gas, v is equivalent to the quantization volume $\Lambda^3 = \left(\frac{h^2}{2\pi m k_B T} \right)^{3/2}$.

From Ω we can obtain the entropy of mixing from $S = k_B \ln \Omega$ with the help of Sterling's approximation $\ln(M!) \approx M \ln(M) - M$:

$$\begin{aligned} S &= k_B (M \ln M - N \ln N - (M - N) \ln(M - N)) \\ &= -M k_B (x \ln x + (1 - x) \ln(1 - x)) \end{aligned} \quad (4)$$

In the last line we introduced a particle fill factor

$$x = N / M$$

which quantifies the fraction of cells that are occupied by particles, and is also known as the mole fraction or the packing ratio. Since $x < 1$, the entropy of mixing is always positive.

For the case of a dilute solution or gas, $N \ll M$, and $(1-x) \approx 1$, so

$$S_{\text{dilute}} \approx -Nk_B \ln x \quad \text{or} \quad -nR \ln x$$

We can derive the ideal gas law $p = Nk_B T / V$ from this result by making use of the thermodynamic identity $p = T (\partial S / \partial V)_N$.

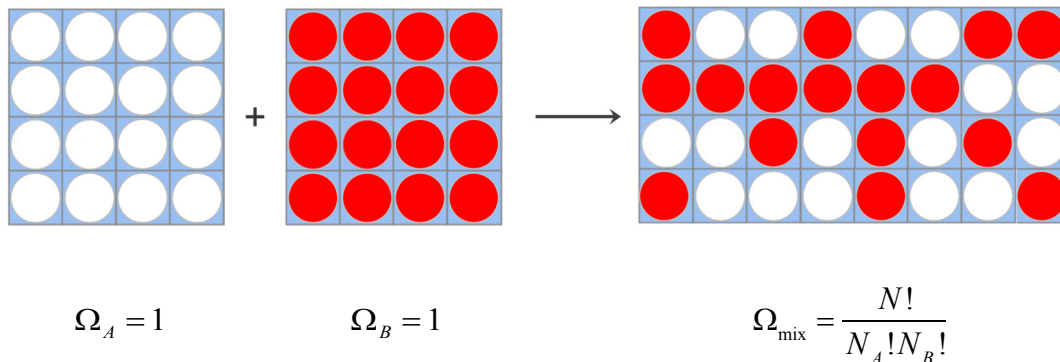
Binary Fluid

Entropy of Mixing

The thermodynamics of the mixing process is important to phase equilibria, hydrophobicity, solubility, and related solvation problems. The process of mixing two pure substances A and B is shown below. We define the composition of the system through the number of A and B particles: N_A and N_B and the total number of particles $N = N_A + N_B$, which also equals the number of cells. We begin with two containers of the homogeneous pure fluids and mix them together, keeping the total number of cells constant. In the case of the pure fluids before mixing, all cells of the container are initially filled, so there is only one accessible microstate, $\Omega_{\text{pure}} = 1$, and

$$S_{\text{pure}} = k_B \ln 1 = 0$$

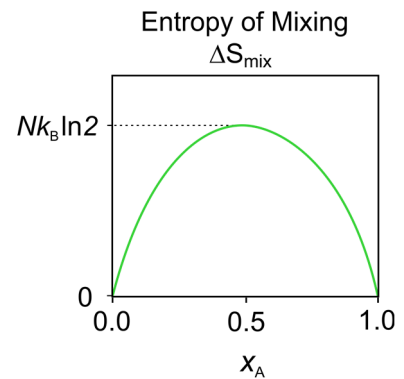
When the two containers are mixed, the number of possible microstates are given by the binomial distribution: $\Omega_{\text{mix}} = N! / N_A! N_B!$.



If these particles have no interactions, each microstate is equally probable, and similar to eq. (4) we obtain the entropy of the mixture as

$$S_{\text{mix}} = -Nk_B (x_A \ln x_A + x_B \ln x_B) \quad (5)$$

For the mixture, we define the mole fractions for the two components: $x_A = N_A / N$ and $x_B = N_B / N$. As before, since x_A and $x_B < 1$, the entropy for the mixture is always positive. The entropy of mixing is then calculated from $\Delta S_{\text{mix}} = S_{\text{mix}} - (S_{\text{pure A}} + S_{\text{pure B}})$. Since the entropy of the pure substances in this model is zero, $\Delta S_{\text{mix}} = S_{\text{mix}}$. A plot of this function as a function of mole fractions illustrates that the maximum entropy mixture has $x_A = x_B = 0.5$.



Ideal systems—be they gasses, solutions, or any variety of molecular ensembles—are characterized by no interactions between particles. Under these conditions, the free energy of mixing is purely entropic with $\Delta A_{\text{mix}} = -T\Delta S_{\text{mix}}$.

Intermolecular Interactions

To look at real (non-ideal) systems, we now add interactions between particles by assigning an interaction energy ω between two cells which are in contact. The interaction energy can be positive (destabilizing) or negative (favorable).



With the addition of intermolecular interactions, each microstate will have a distinct energy, the canonical partition function can be obtained from eq. (1), and other thermodynamic properties follow.

In the case of a mixture, we assign separate interaction energies for each adjoining A-A, B-B, or A-B pair in a given microstate: $\omega_{AA}, \omega_{BB}, \omega_{AB}$. How do we calculate the energy of a microstate? m is the total number of molecular contacts in the volume, and these can be divided into A-A, B-B, or A-B contacts:

$$m = m_{AA} + m_{BB} + m_{AB}$$

While m is constant, the counts of specific contacts m_{ij} vary by microstate. Then the energy of the mixture *for the single i^{th} microstate* can be written as

$$E_{\text{mix}} = m_{AA}\omega_{AA} + m_{BB}\omega_{BB} + m_{AB}\omega_{AB} \quad (6)$$

and the internal energy comes from an ensemble average of this quantity. An exact calculation of the internal energy from the partition function would require a sum over all possible configurations with their individual contact numbers. Instead, we can use a simpler, approximate approach which uses a strategy that starts by expressing each term in eq. (6) in terms of m_{AB} . We know:

$$\begin{aligned} m_{AA} &= (\text{Total contacts for A}) - (\text{Contacts of A with B}) \\ &= \frac{zN_A}{2} - \frac{m_{AB}}{2} \end{aligned} \quad (7)$$

$$m_{BB} = \frac{zN_B}{2} - \frac{m_{AB}}{2} \quad (8)$$

Then we have

$$\begin{aligned}
E_{\text{mix}} &= \left(\frac{z\omega_{AA}N_A}{2} \right) + \left(\frac{z\omega_{BB}N_B}{2} \right) + m_{AB} \left(\omega_{AB} - \frac{\omega_{AA} + \omega_{BB}}{2} \right) \\
&= U_{\text{pure } A} + U_{\text{pure } B} + m_{AB} \left(\omega_{AB} - \frac{\omega_{AA} + \omega_{BB}}{2} \right)
\end{aligned} \tag{9}$$

Here in the second step, we recognize that the first two terms are just the energy of the two pure liquids before mixing. These are calculated by taking the number of cells in the pure liquid (N_i) times the number of contacts per cell (z) and then divide by two, so you do not double count the contacts.

$$U_{\text{pure},i} = \frac{z\omega_{ii}N_i}{2} \tag{10}$$

Equation (9) describes the energy of a microstate in terms of m_{AB} . To simplify our calculation of U_{mix} , we make a “mean field approximation,” which replaces m_{AB} with its statistical average $\langle m_{AB} \rangle$:

$$\begin{aligned}
\langle m_{AB} \rangle &= (\# \text{ of contact sites for A}) \times (\text{probability of contact site being B}) \\
&= (N_A z) \left(\frac{N_B}{N} \right) = z x_A x_B N
\end{aligned} \tag{11}$$

Then for the energy for the mixed state $U_{\text{mix}} = \langle E_{\text{mix}} \rangle$, we obtain:

$$U_{\text{mix}} = U_{\text{pure } A} + U_{\text{pure } B} + x_A x_B N k_B T \chi_{AB} \tag{12}$$

Here we have introduced the unitless *exchange parameter*,

$$\chi_{AB} = \frac{z}{k_B T} \left(\omega_{AB} - \frac{\omega_{AA} + \omega_{BB}}{2} \right) = \frac{z \Delta\omega}{k_B T} \tag{13}$$

which measures the relative change of intermolecular interaction for one cell in the lattice switching from an A-A and B-B contact to 2 A-B contacts. This average change of interaction energies is expressed in units of kBT. Dividing by z gives the average interaction energy per contact.

$$\chi_{AB} > 0 \rightarrow \text{unfavorable A}\cdot\text{B interaction}$$

$$\chi_{AB} < 0 \rightarrow \text{favorable A}\cdot\text{B interaction}$$

We can now determine the change in internal energy on mixing:

$$\begin{aligned}
\Delta U_{\text{mix}} &= (U_{\text{mix}} - U_{\text{pure } A} - U_{\text{pure } B}) \\
&= x_A x_B N k_B T \chi_{AB}
\end{aligned} \tag{14}$$

Note ΔU_{mix} as a function of composition has its maximum value for a mixture with $x_A=0.5$.

Note that in the mean field approximation, the canonical partition function is

$$Q = \frac{N!}{N_A!N_B!} q_A^{N_A} q_B^{N_B} \exp[-U_{\text{mix}} / k_B T]$$

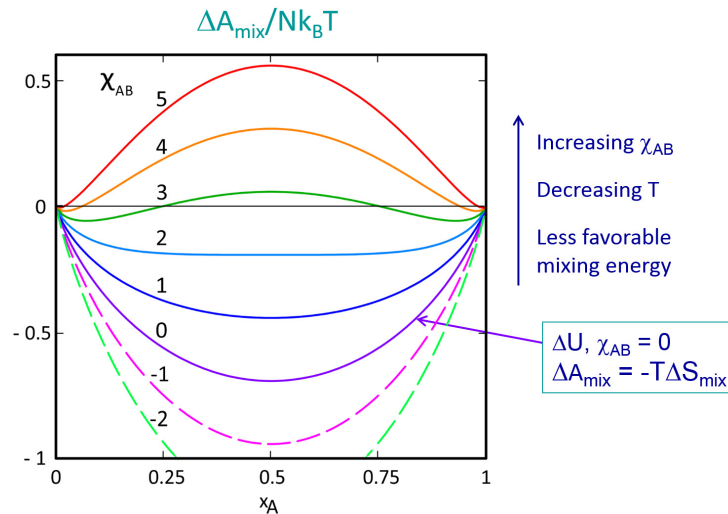
We kept the internal molecular partition functions here for completeness, but for the simple particles in this model $q_A = q_B = 1$.

Free Energy of Mixing¹

Using eqs. (5) and (14), we can now obtain the free energy of mixing:

$$\begin{aligned} \Delta A_{\text{mix}} &= \Delta U_{\text{mix}} - T \Delta S_{\text{mix}} \\ &= N k_B T (x_A x_B \chi_{AB} + x_A \ln x_A + x_B \ln x_B) \end{aligned}$$

This function is plotted below as a function of mole fraction for different values of the exchange parameter. When there are no intermolecular interactions ($\chi_{AB} = 0$), the mixing is spontaneous for any mole fraction and purely entropic. Any strongly favorable A-B interaction ($\chi_{AB} < 0$) only serves to decrease the free energy further for all mole fractions.

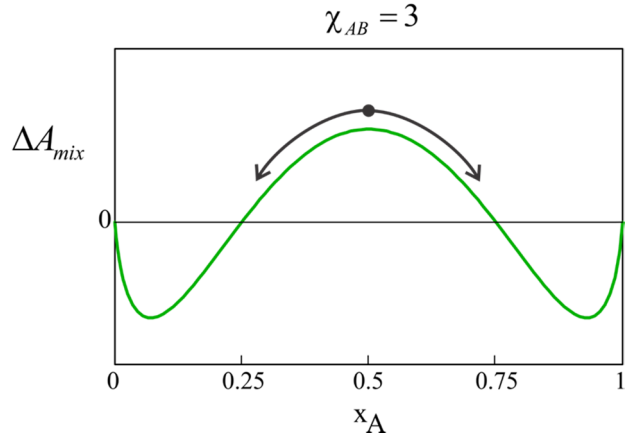


As χ_{AB} increases, we see the free energy for mixing rise, with the biggest changes for the 50/50 mixture. To describe the consequences, let's look at the curve for $\chi_{AB} = 3$, for which certain compositions are miscible ($\Delta A_{\text{mix}} < 0$) and others immiscible ($\Delta A_{\text{mix}} > 0$).

Consider what would happen if we prepare a 50/50 mixture of this solution. The free energy of mixing is positive at the equilibrium composition of the $x_A = 0.5$ homogeneous mixture, indicating that the two components are immiscible. However, there are other mixture compositions that do have a negative free energy of mixing. Under these conditions the solution can separate into two

1. J. H. Hildebrand and R. L. Scott, *Regular Solutions*. (Prentice-Hall, Englewood Cliffs, N.J., 1962).

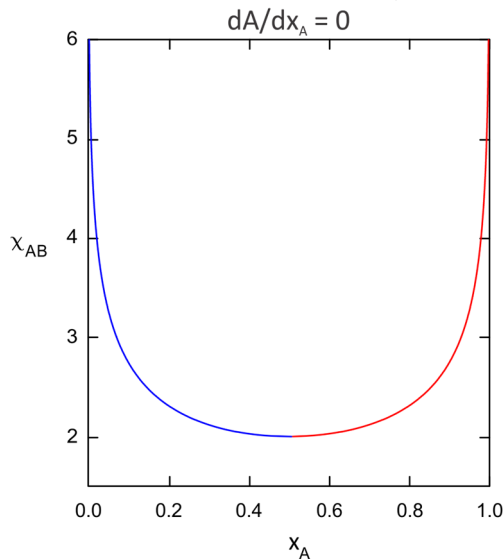
phases in such a way that ΔA_{mix} is minimized. This occurs at mole fractions of $x_A = 0.07$ & 0.93 , which shows us that one phase will be characterized by $x_A \gg x_B$ and the other with $x_A \ll x_B$. If we prepare an unequal mixture with positive ΔA_{mix} , for example $x_A = 0.3$, the system will still spontaneously phase separate although mass conservation will dictate that the total mass of the fraction with $x_A = 0.07$ will be greater than the mass of the fraction at $x_A = 0.93$. As χ_{AB} increases beyond 3, the mole fraction of the lesser component decreases as expected for the hydrophobic effect. Consider if $A = \text{water}$ and $B = \text{oil}$. ω_{BB} and ω_{AB} are small and negative, ω_{AA} is large and negative, and $\chi_{AB} \gg 1$.



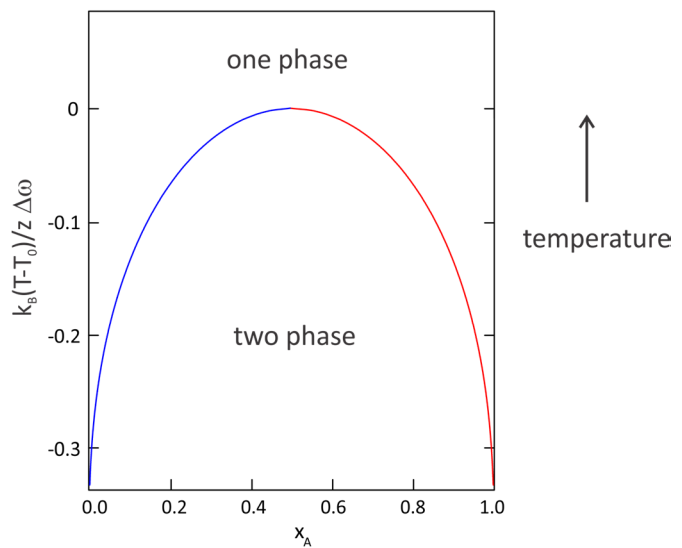
Critical Behavior

Note that 50/50 mixtures with $2 < \chi_{AB} < 2.8$ have a negative free energy of mixing to create a single homogeneous phase, yet, the system can still lower the free energy further by phase separating. As seen in the figure, $\chi_{AB} = 2$ marks a crossover from one phase mixtures to two phase mixtures, which is the signature of a critical point. We can find the conditions for phase equilibria by locating the free energy minima as a function of χ_{AB} , which leads to the phase diagrams as a function of χ_{AB} and T below. The critical temperature for crossover from one- to two-phase behavior is T_0 , and $\Delta\omega$ is the average differential change in interaction energy defined in eq. (13).

Conditions for phase equilibria in binary fluid mixture



T-dependence of phase equilibria in binary fluid mixture



Readings

1. K. Dill and S. Bromberg, *Molecular Driving Forces: Statistical Thermodynamics in Biology, Chemistry, Physics, and Nanoscience*. (Taylor & Francis Group, New York, 2010).
2. W. W. Graessley, *Polymeric Liquids and Networks: Structure and Properties*. (Garland Science, New York, 2004), Ch. 3.